

Higher-Order Results in the Electroweak Theory

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1. Introduction
2. Present status of electroweak precision observables
3. Remaining theoretical uncertainties
4. Conclusions

1. Introduction

Electroweak precision measurements:

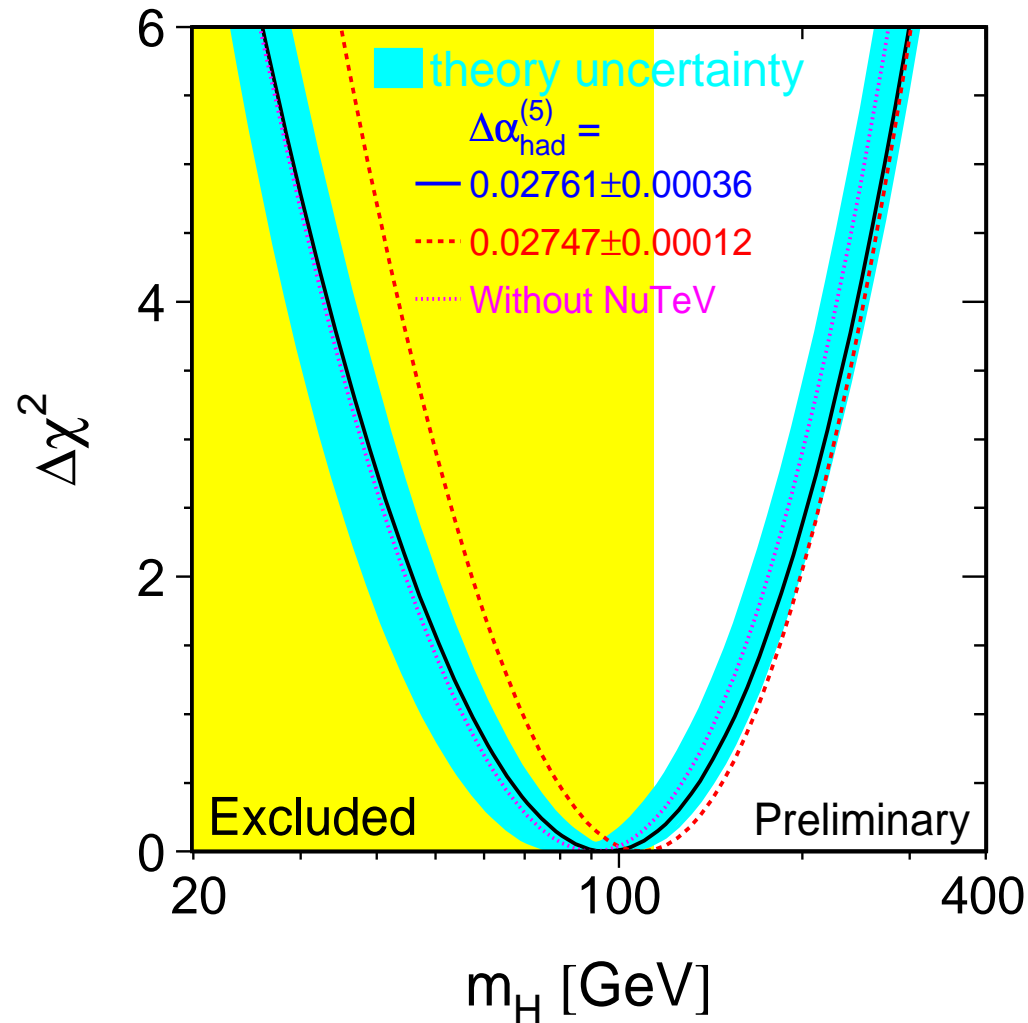
M_Z [GeV]	=	91.1875 ± 0.0021	0.002%
G_μ [GeV ⁻²]	=	$1.16637(1) 10^{-5}$	0.0009%
m_t [GeV]	=	174.3 ± 5.1	2.9%
M_W [GeV]	=	80.426 ± 0.034	0.04%
$\sin^2 \theta_{\text{eff}}^{\text{lept}}$	=	0.23150 ± 0.00016	0.07%
Γ_Z [GeV]	=	2.4952 ± 0.0023	0.09%
...			

Quantum effects of the theory:

Loop corrections to “pseudo-observables” $M_W, \sin^2 \theta_{\text{eff}}, \Gamma_Z, \dots$: $\sim \mathcal{O}(1\%)$

\Rightarrow Indirect constraints on M_H, \dots — effects of “new physics”?

Global fit to all data in the SM: [LEPEWWG '03]



Theoretical uncertainties:

– unknown higher-order corrections

⇒ “blue band”

– exp. error of input parameters:

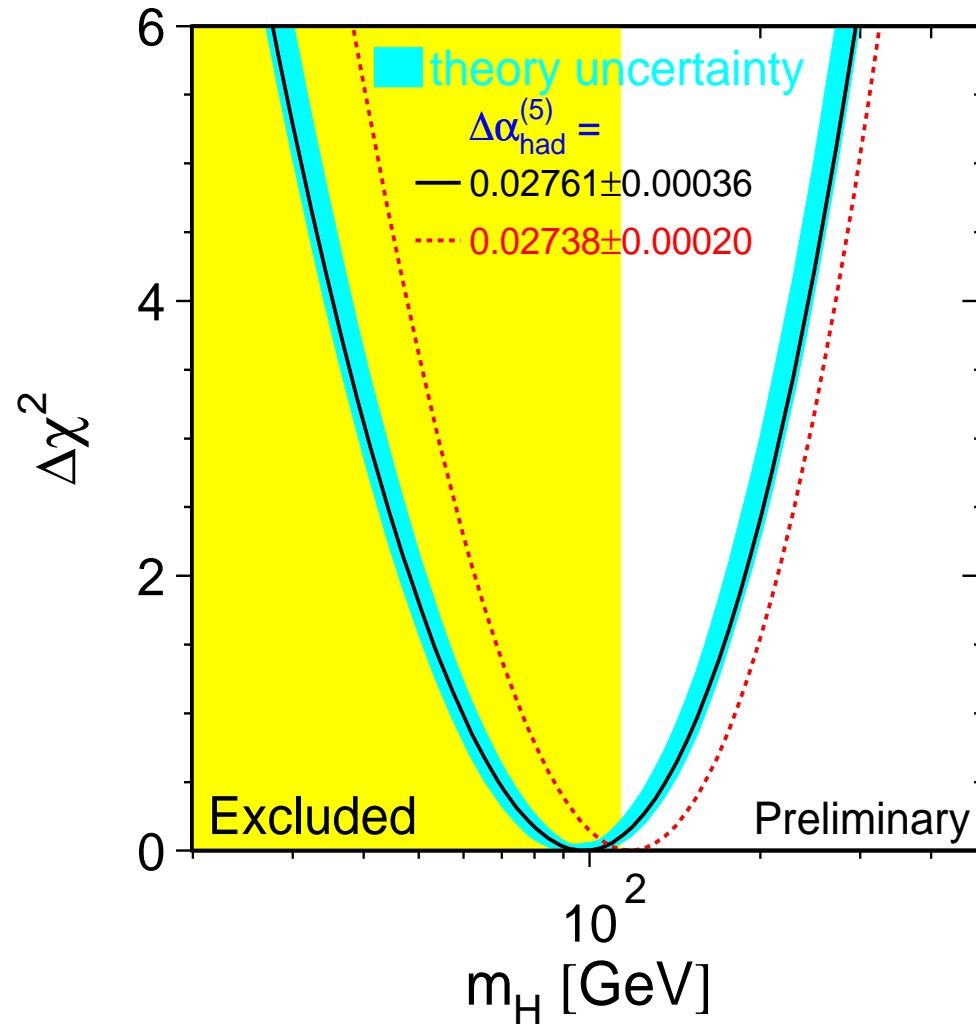
$m_t, \Delta\alpha_{\text{had}}, \dots$

Overall fit probability (quality of the fit): 4.4% (27%) with (without) NuTeV

⇒ $M_H = 91_{-37}^{+58}$ GeV, $M_H < 211$ GeV, 95% C.L.

Uncertainties from unknown higher-order corrections:

Comparison: global fit results Winter '01 [*LEPEWWG '01*]



$$\Rightarrow M_H = 98_{-38}^{+58} \text{ GeV}$$

$$M_H < 212 \text{ GeV, 95\% C.L.}$$

⇒ “Blue band” has widened

2. Present status of electroweak precision observables

Theoretical predictions for M_W , $\sin^2 \theta_{\text{eff}}$, ... in the SM:

Definition of Fermi constant G_μ via muon lifetime:

$$\tau_\mu^{-1} = \frac{G_\mu^2 m_\mu^5}{192\pi^3} F\left(\frac{m_e^2}{m_\mu^2}\right) \left(1 + \frac{3}{5} \frac{m_\mu^2}{M_W^2}\right) (1 + \Delta q)$$

Δq : QED corrections in Fermi Model [*R. Behrends, R. Finkelstein, A. Sirlin '56*]
[*T. van Ritbergen, R. Stuart '99*] [*M. Steinhauser, T. Seidensticker '99*]

Comparison with SM prediction for muon decay:

$$\Rightarrow M_W^2 \left(1 - \frac{M_W^2}{M_Z^2}\right) = \frac{\pi\alpha}{\sqrt{2}G_\mu} (1 + \Delta r),$$



loop corrections

\Rightarrow Theoretical prediction for M_W in terms of M_Z , α , G_μ , $\Delta r(m_t, M_H, \dots)$

Effective couplings at the Z resonance:

$$\Rightarrow \sin^2 \theta_{\text{eff}} = \frac{1}{4} \left(1 - \frac{\text{Re } g_V}{\text{Re } g_A} \right) = \left(1 - \frac{M_W^2}{M_Z^2} \right) \kappa(M_W^2)$$

One-loop result for M_W in the SM:

[A. Sirlin '80], [W. Marciano, A. Sirlin '80]

$$\Delta r_{1\text{-loop}} = \Delta\alpha - \frac{c_W^2}{s_W^2} \Delta\rho + \Delta r_{\text{rem}}(M_H)$$

$$\sim \log \frac{M_Z}{m_f} \quad \sim m_t^2$$

$$\sim 6\% \quad \sim 3.3\% \quad \sim 1\%$$

Leading contributions to M_W , $\sin^2 \theta_{\text{eff}}$, ... from mass splitting between isospin doublet fields enter via

$$\Delta\rho = \frac{\Sigma_Z(0)}{M_Z^2} - \frac{\Sigma_W(0)}{M_W^2}, \quad \Delta M_W \approx \frac{M_W}{2} \frac{c_W^2}{c_W^2 - s_W^2} \Delta\rho, \quad \Delta \sin^2 \theta_{\text{eff}} \approx - \frac{c_W^2 s_W^2}{c_W^2 - s_W^2} \Delta\rho$$

Two-loop QCD corrections:

[A. Djouadi, C. Verzegnassi '87], [A. Djouadi '89], [B. Kniehl '90], [F. Halzen, B. Kniehl '91]

Electroweak two-loop corrections:

$\sin^2 \theta_{\text{eff}}$:

Expansion for asymptotically large m_t :

Leading m_t^4/M_W^4 terms (enter via $\Delta\rho$)

[J. van der Bij, F. Hoogeveen '87], [R. Barbieri, M. Beccaria, P. Ciafaloni, G. Curci, A. Vicere '92], [J. Fleischer, O.V. Tarasov, F. Jegerlehner '93]

Next-to-leading m_t^2/M_W^2 terms [G. Degrassi, P. Gambino, A. Sirlin '97]

Expansions: $M_W, M_Z, M_H \ll m_t$, $M_W, M_Z \ll m_t, M_H$ + interpolation

→ Same order of magnitude as m_t^4/M_W^4 terms, same sign

Expansion for asymptotically large M_H :

Leading M_H^2/M_W^2 terms (enter via $\Delta\rho$)

[J. van der Bij, M. Veltman '84], [J. van der Bij '84]

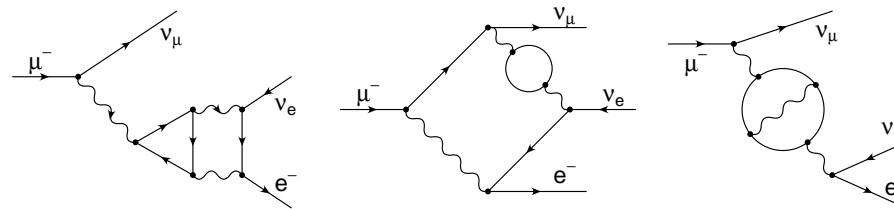
M_W :

Complete electroweak two-loop result:

Fermionic two-loop contributions:

[A. Freitas, W. Hollik, W. Walter, G.W. '00, '02]

[M. Awramik, M. Czakon '03]



⇒ full dependence on m_t , complete light fermion contributions

Method:

Extraction of QED contributions in the Fermi model

Tensor integral reduction for general two-loop 2-point functions

[G. W. '92], [G. W., R. Scharf, M. Böhm '94]

General R_ξ gauge, careful treatment of γ_5 in dimensional regularisation

[*G. 't Hooft, M. Veltman '72*], [*P. Breitenlohner, D. Maison '77*]

Two-loop on-shell renormalization (physical parameters)

renormalization in gauge-fixing and ghost sector

mass renormalization of unstable particles

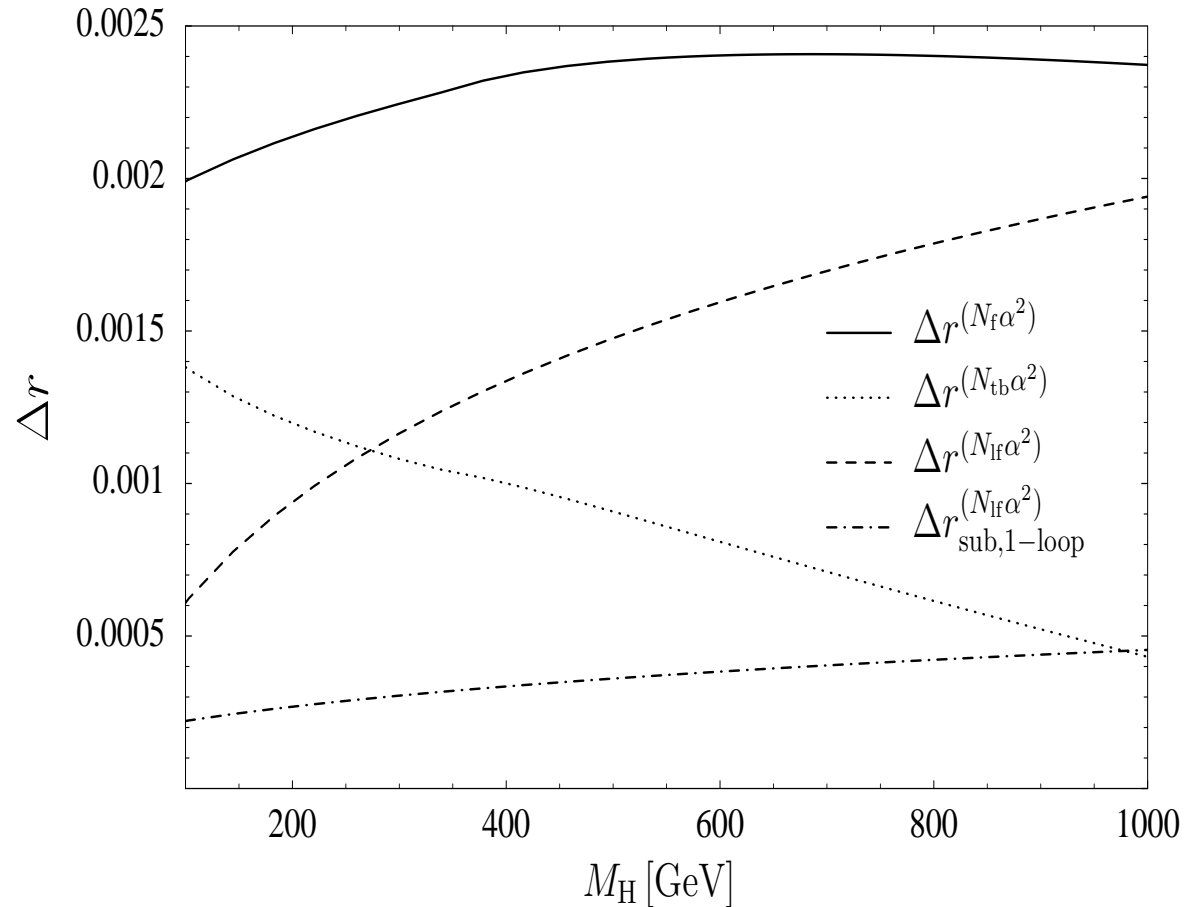
⇒ gauge-invariant definition according to real part of complex pole

Numerical evaluation via one-dimensional integral representations with elementary functions, results for arbitrary mass values

[*S. Bauberger, F. Berends, M. Böhm, M. Buza '95*], [*S. Bauberger, M. Böhm '95*]

Comparison of contributions from top/bottom doublet and from light fermions:

[A. Freitas, W. Hollik, W. Walter, G.W. '02]

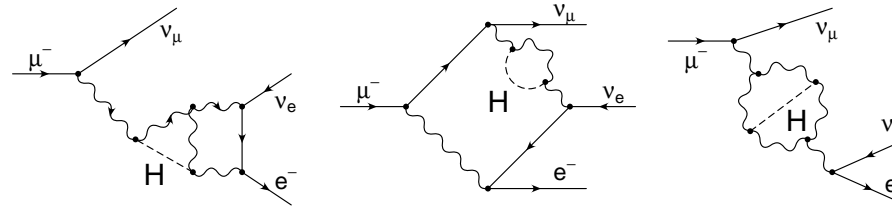


$$\Delta r_{\text{sub},1\text{-loop}}^{(N_{lf} \alpha^2)} = \Delta r^{(N_{lf} \alpha^2)} - 2\Delta\alpha\Delta r_{\text{bos}}^{(\alpha)}$$

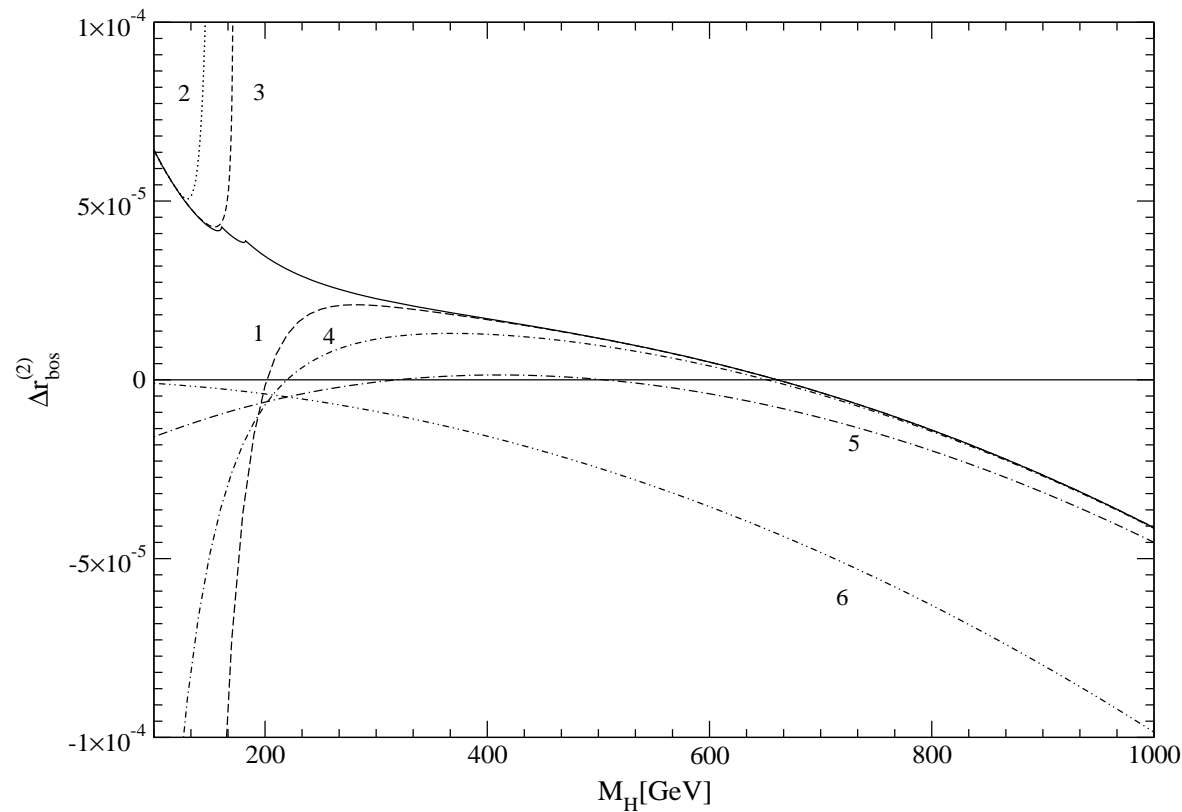
⇒ non-leading t - b and genuine 2-loop light fermion contributions lead to shift in M_W of about 5 MeV

Purely bosonic two-loop contributions:

[M. Awramik, M. Czakon '02] [A. Onishchenko, O. Veretin '02]



Similar methods as for fermionic contributions, numerical integration & expansion in $(M_H^2 - M_Z^2)$, $(M_Z^2 - M_W^2)$ and for $M_H \gg M_Z$



$$\Rightarrow \Delta M_W \lesssim \pm 1 \text{ MeV}$$

Three-loop QCD corrections to M_W and $\sin^2 \theta_{\text{eff}}$:

[*L. Avdeev, J. Fleischer, S.M. Mikhailov, O. Tarasov '94*]

[*K. Chetyrkin, J. Kühn, M. Steinhauser '95*]

Leading EW corrections beyond 2-loop:

Pure fermion-loop contributions up to 4-loop order

[*A. Stremplatt '98*], [*G.W. '98*]

Leading 3-loop corrections to $\Delta\rho$:

$\mathcal{O}(G_\mu^3 m_t^6)$, $\mathcal{O}(G_\mu^2 \alpha_s m_t^4)$ terms, limit: $M_H = 0$

[*J. van der Bij, K. Chetyrkin, M. Faisst, G. Jikia, T. Seidensticker '01*]

$\mathcal{O}(G_\mu^3 m_t^6)$, $\mathcal{O}(G_\mu^2 \alpha_s m_t^4)$ terms, arbitrary M_H

[*M. Faisst, J. Kühn, T. Seidensticker, O. Veretin '03*]

Method:

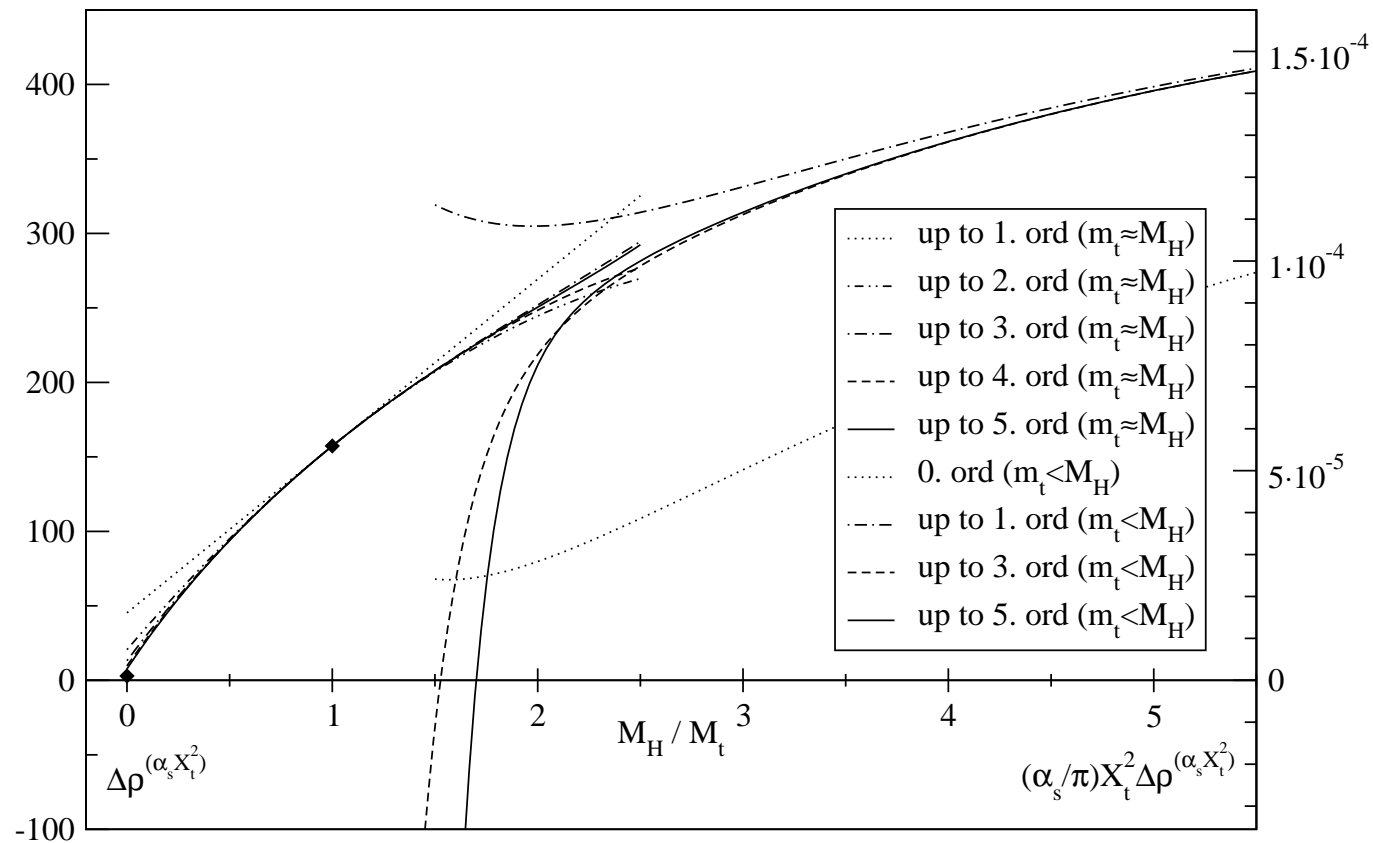
Expansions around $M_H = m_t$ and for $M_H \gg m_t$

\Rightarrow three-loop vacuum integrals with one mass scale

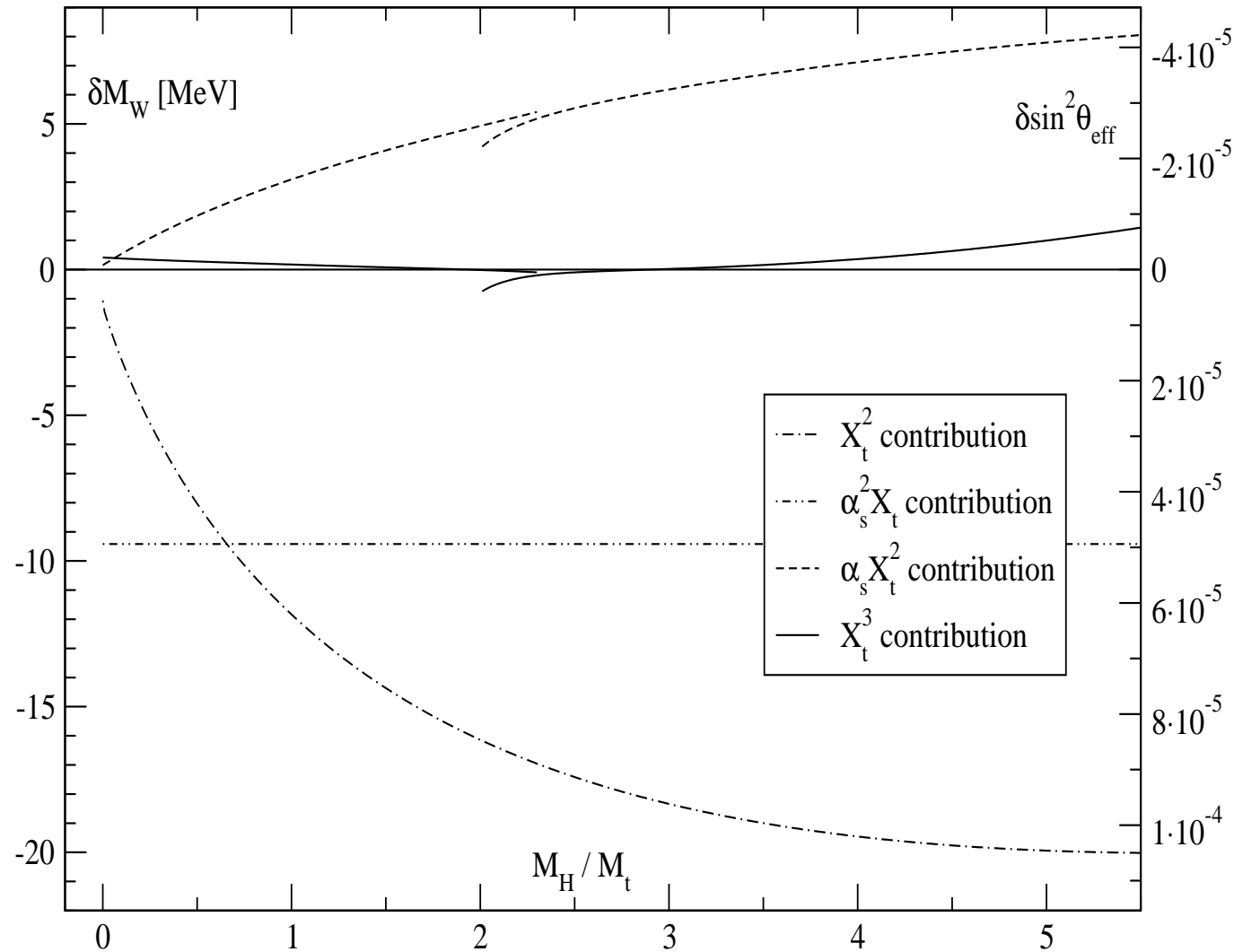
Transition between \overline{MS} and on-shell value of top-quark mass

⇒ two-loop top-quark self-energy on-shell, evaluated with expansions around $M_H = m_t$ and for $M_H \gg m_t$

$\mathcal{O}(G_\mu^2 \alpha_s m_t^4)$ contributions to $\Delta\rho$:



Shifts in M_W , $\sin^2 \theta_{\text{eff}}$ from $\mathcal{O}(G_\mu^3 m_t^6)$, $\mathcal{O}(G_\mu^2 \alpha_s m_t^4)$ contributions to $\Delta\rho$
 [M. Faisst, J. Kühn, T. Seidensticker, O. Veretin '03]



$\Rightarrow \Delta M_W \lesssim 5 \text{ MeV}, \Delta \sin^2 \theta_{\text{eff}} \lesssim -3 \times 10^{-5}$

Simple parameterization of full result for M_W (contains all known corrections): [M. Awramik, M. Czakon, A. Freitas, G.W. '03]

$$\begin{aligned} M_W &= M_W^0 - c_1 dH - c_2 dH^2 + c_3 dH^4 + c_4 dh \\ &\quad - c_5 d\alpha + c_6 dt - c_7 dt^2 - c_8 dH dt \\ &\quad + c_9 dh dt - c_{10} d\alpha_s + c_{11} dz \end{aligned}$$

where

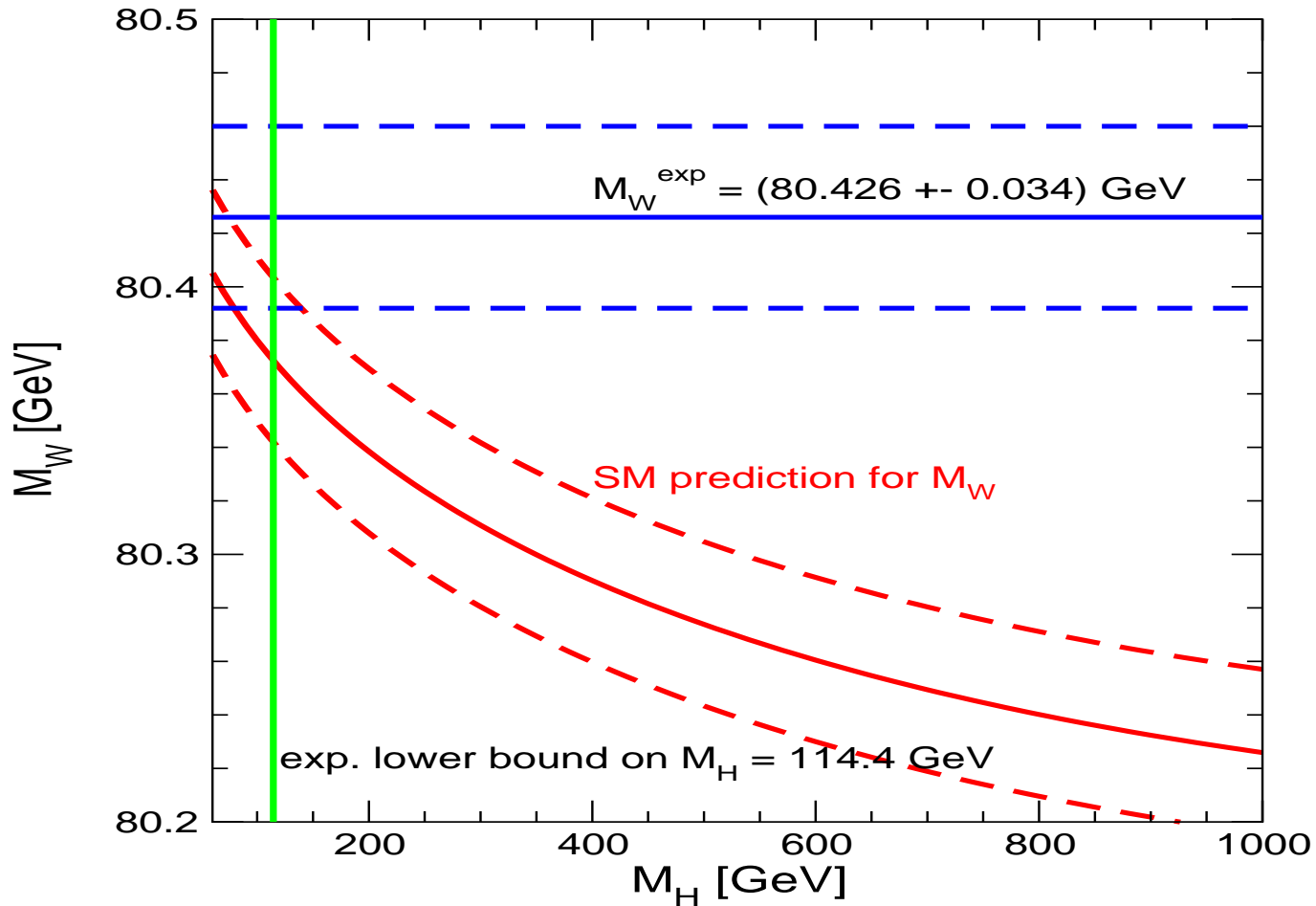
$$dH = \ln \left(\frac{M_H}{100 \text{ GeV}} \right), \quad dh = \left(\frac{M_H}{100 \text{ GeV}} \right)^2 - 1, \quad dt = \left(\frac{m_t}{174.3 \text{ GeV}} \right)^2 - 1$$

$$d\alpha = \frac{\Delta\alpha}{0.05907} - 1, \quad d\alpha_s = \frac{\alpha_s(M_Z)}{0.119} - 1, \quad dz = M_Z - 91.1875 \text{ GeV}$$

\Rightarrow approximates full result within 0.5 MeV for $10 \text{ GeV} \leq M_H \leq 1 \text{ TeV}$,
2 σ variations

SM prediction for M_W vs. experimental result:

[M. Awramik, M. Czakon, A. Freitas, G.W. '03]



\Rightarrow Light Higgs preferred

$M_W^{\text{exp}} = 80.426 \pm 0.034 \text{ GeV}$, $m_t^{\text{exp}} = 174.3 \pm 5.1 \text{ GeV}$,

$M_H > 114.4 \text{ GeV}$, 95% C.L.

$\Rightarrow 1\sigma$ regions hardly overlap

3. Remaining theoretical uncertainties

Uncertainties from unknown higher-order corrections, different methods:

“Traditional Blue Band method”:

[*D. Bardin, M. Grünewald, G. Passarino '99*]

Size of unknown corrections estimated by varying different options in

ZFITTER [*D. Bardin et al.*], *TOPAZ0* [*G. Passarino*]

Different renormalization schemes, different resummations of non-leading terms

⇒ differ by higher-order corrections

difficult to quantify, genuine effects of irreducible higher-order corrections not accessible

Error estimate for blue band plot up to Moriond '01: [*LEPEWWG '01*]

$$\Delta M_W^{\text{theo}} \approx \pm 3 \text{ MeV}, \quad \Delta \sin^2 \theta_{\text{eff}}^{\text{theo}} \approx \pm 3 \times 10^{-5}$$

More conservative estimate for $\sin^2 \theta_{\text{eff}}$ (since summer '01) [LEPEWWG '01]

Effect of non-leading fermionic and purely bosonic two-loop corrections on $\sin^2 \theta_{\text{eff}}$ not yet known

⇒ use corresponding shift in M_W to estimate uncertainty of $\sin^2 \theta_{\text{eff}}$

$$\text{according to } \sin^2 \theta_{\text{eff}} = \left(1 - \frac{M_W^2}{M_Z^2}\right) \kappa(M_W^2)$$

⇒ upward shift in $\sin^2 \theta_{\text{eff}}$ of $\Delta \sin^2 \theta_{\text{eff}}^{\text{theo}} \approx 8 \times 10^{-5}$ for $M_H = 100$ GeV

Estimate of QCD corrections:

[A. Freitas, W. Hollik, W. Walter, G.W. '00, '02]

– on-shell vs. $\overline{\text{MS}}$ value of m_t

– estimate $\Delta_r(\alpha^2 \alpha_s) / \Delta_r(\alpha^2)$ from known $\Delta_r(\alpha \alpha_s) / \Delta_r(\alpha)$

⇒ $\mathcal{O}(\alpha^2 \alpha_s)$: $\Delta M_W^{\text{theo}} \approx 4$ MeV, $\mathcal{O}(\alpha \alpha_s^3)$: $\Delta M_W^{\text{theo}} \approx 1$ MeV

Transition from fixed-width definition to running width definition of M_W :

[A. Freitas, W. Hollik, W. Walter, G.W. '00, '02]

$$\Delta M_W^{\text{theo}} \approx 1 \text{ MeV}$$

Estimate according to possible enhancement factors multiplied by coefficients of $\mathcal{O}(1)$

[U. Baur, R. Clare, J. Erler, S. Heinemeyer, D. Wackeroth, G. W., D. Wood '01]

order	sector	estimate	M_W [MeV]	$\sin^2 \theta_{\text{eff}}$ (10^5)
α^2	fermionic	$N(\alpha/4\pi\hat{s}^2)^2$	complete	4.1
α^2	bosonic	$(\alpha/\pi\hat{s}^2)^2$	2.1	4.1
$\alpha\alpha_s^2$	t - b doublet	$N_C C_F C_A \alpha\alpha_s^2 / 4\pi^3 \hat{s}^2$	complete	complete
$\alpha\alpha_s^2$	light doubl.	$2 N_C C_F C_A \alpha\alpha_s^2 / 4\pi^3 \hat{s}^2$	complete	complete
$\alpha^3 m_t^6$	heavy top	$5.3 N_C^2 (\alpha m_t^2 / 4\pi\hat{s}^2 M_W^2)^3$	4.1	2.3
$\alpha^3 m_t^6$	heavy top	$3.3 N_C (\alpha m_t^2 / 4\pi\hat{s}^2 M_W^2)^3$	0.9	0.5
$\alpha^2 \alpha_s m_t^4$	heavy top	$3.9 N_C C_F \alpha^2 \alpha_s m_t^4 / 16\pi^3 \hat{s}^4 M_W^4$	4.5	2.5
$\alpha\alpha_s^3 m_t^2$	heavy top	$N_C C_F C_A^2 \alpha\alpha_s^3 m_t^2 / 4\pi^4 \hat{s}^2 M_W^2$	1.3	0.8
	total		7	7

⇒ good agreement for the corrections that meanwhile have been calculated

Estimate of remaining uncertainties from unknown higher orders in M_W and $\sin^2 \theta_{\text{eff}}$:

[M. Awramik, M. Czakon, A. Freitas, G.W. '03]

M_W : $\mathcal{O}(G_\mu^3 m_t^4 M_Z^2)$, $\mathcal{O}(G_\mu^2 \alpha_s m_t^2 M_Z^2)$, $\mathcal{O}(\alpha \alpha_S^3)$

$$\Delta M_W^{\text{theo}} \approx 4 \text{ MeV}$$

$\sin^2 \theta_{\text{eff}}$: non-leading fermionic two-loop contributions, purely bosonic two-loop contributions, $\mathcal{O}(G_\mu^3 m_t^4 M_Z^2)$, $\mathcal{O}(G_\mu^2 \alpha_s m_t^2 M_Z^2)$, $\mathcal{O}(\alpha \alpha_S^3)$

$$\Delta \sin^2 \theta_{\text{eff}}^{\text{theo}} \approx 6 \times 10^{-5}$$

Uncertainties from experimental errors of input parameters:

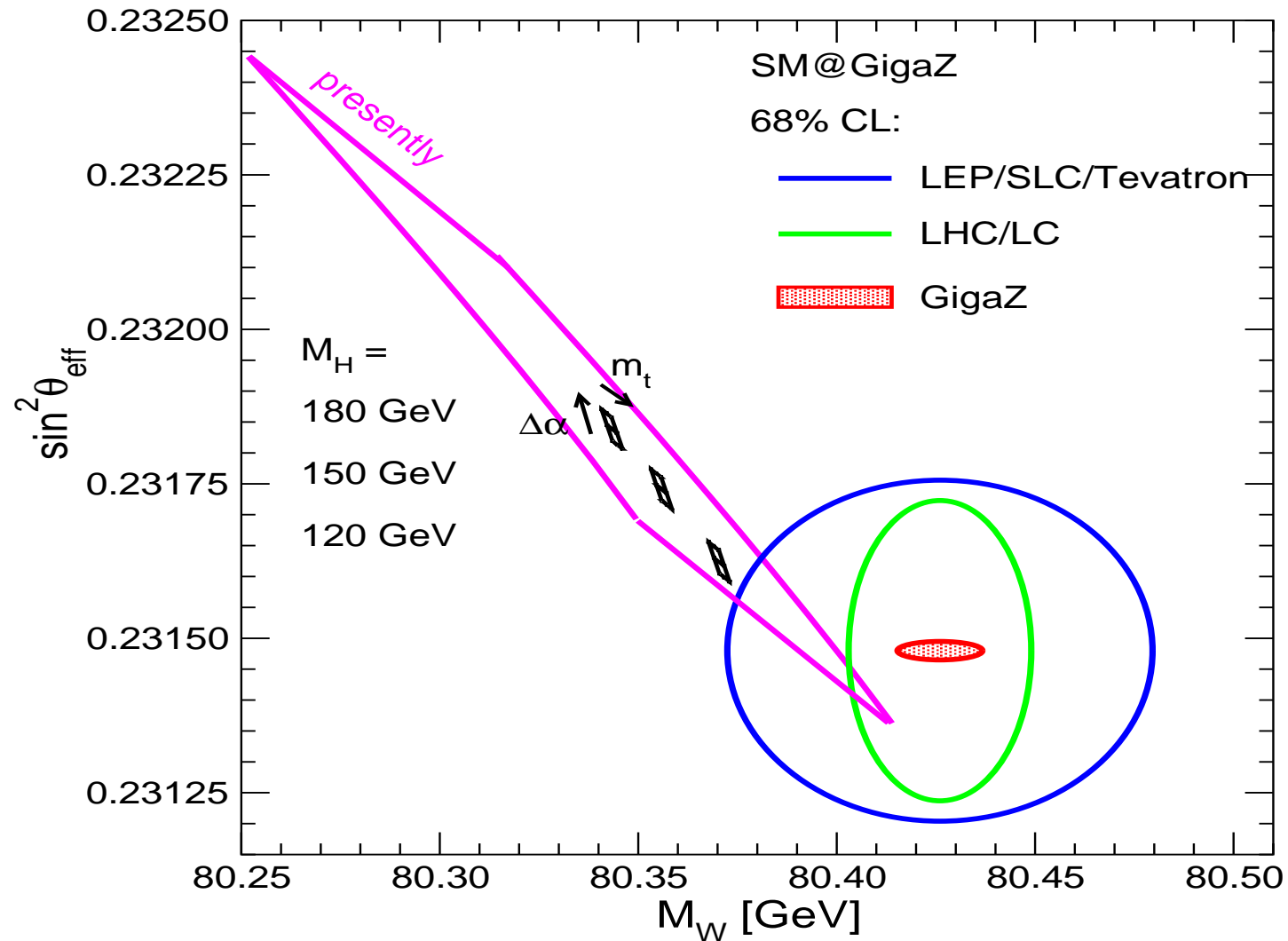
$$\begin{aligned}
 \Delta m_t = \pm 5.1 \text{ GeV} &\Rightarrow \Delta M_W \approx \pm 31 \text{ MeV}, & \Delta \sin^2 \theta_{\text{eff}} &\approx \pm 16 \times 10^{-5} \\
 \Delta(\Delta\alpha_{\text{had}}) = \pm 3.6 \times 10^{-4} &\Rightarrow \Delta M_W \approx \pm 6.5 \text{ MeV}, & \Delta \sin^2 \theta_{\text{eff}} &\approx \pm 13 \times 10^{-5} \\
 \Delta\alpha_s = \pm 0.002 &\Rightarrow \Delta M_W \approx \pm 1.5 \text{ MeV}, & \Delta \sin^2 \theta_{\text{eff}} &\approx \pm 2.5 \times 10^{-5} \\
 \Delta M_Z = \pm 2.1 \text{ MeV} &\Rightarrow \Delta M_W \approx \pm 2.5 \text{ MeV}, & \Delta \sin^2 \theta_{\text{eff}} &\approx \pm 1.4 \times 10^{-5}
 \end{aligned}$$

Possible future improvement in $\Delta(\Delta\alpha_{\text{had}})$: $\rightarrow \pm 0.5 \times 10^{-4}$ [*F. Jegerlehner '01*]

Prospects for experimental accuracies at the next generation of colliders:

$$\begin{aligned}
 \Delta m_t^{\text{exp}} / \text{GeV} : & \quad 5.1 \xrightarrow{\text{Tev/LHC}} 1.5 \xrightarrow{\text{LC}} 0.1 \\
 \Delta M_W^{\text{exp}} / \text{MeV} : & \quad 34 \xrightarrow{\text{Tev/LHC}} 15 \xrightarrow{\text{LC}} 7 \\
 \Delta \sin^2 \theta_{\text{eff}}^{\text{exp}} / 10^{-5} : & \quad 16 \xrightarrow{\text{Tev/LHC}} 16 \xrightarrow{\text{LC}} 1
 \end{aligned}$$

SM prediction for M_W and $\sin^2 \theta_{\text{eff}}$ vs. current experimental result (LEP2/Tevatron) and prospective accuracies at the LHC and a LC with low-energy option (GigaZ): [J. Erler, S. Heinemeyer, W. Hollik, G.W., P. Zerwas '00]



⇒ Sensitive test of the electroweak theory

4. Conclusions

- Higher-order results for M_W , $\sin^2 \theta_{\text{eff}}$ in the SM:
 M_W : complete two-loop result, leading QCD and ew three-loop corr.
 $\sin^2 \theta_{\text{eff}}$: non-leading two-loop contributions still missing
- Estimate of remaining theoretical uncertainties from unknown higher-order corrections:
 $\Delta M_W^{\text{theo}} \approx 4 \text{ MeV}$, $\Delta \sin^2 \theta_{\text{eff}}^{\text{theo}} \approx 6 \times 10^{-5}$
- Effect of parametric uncertainty of m_t :
present: $\Delta M_W \approx \pm 31 \text{ MeV}$, $\Delta \sin^2 \theta_{\text{eff}} \approx \pm 16 \times 10^{-5}$
LC: $\Delta M_W \approx \pm 1 \text{ MeV}$, $\Delta \sin^2 \theta_{\text{eff}} \approx \pm 0.3 \times 10^{-5}$
- At next generation of colliders:
Improved accuracy of precision observables M_W , $\sin^2 \theta_{\text{eff}}$, m_h , ...
and input parameters m_t , $m_{\tilde{t}}$, ...
 \Rightarrow Very sensitive test of electroweak theory