

Effects of New Physics on CP violation in B decays

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Outline

- $b \leftrightarrow s$ Phenomena: Current Frontier for NP
- Flavor Models w/ *1 Extra CP Phase* in $b \leftrightarrow s$
 - 4th Generation
 - Large s_R - b_R Mixing w/ SUSY: Light $\tilde{s}b_{1R}$ Squark
- 4th Generation & $b \rightarrow s \ell^+ \ell^-$, Δm_{B_S} , $\sin 2\Phi_{B_S}$
 - Does Not Touch EM/Strong Penguins Much
- Light $\tilde{s}b_{1R}$ and $S_{\phi K_S}$, $S_{\eta' K_S}$, Δm_{B_S} , $\sin 2\Phi_{B_S}$
 - Survive $b \rightarrow s \gamma$
 - Tune for $S_{\phi K_S}$, $S_{\eta' K_S}$ Possible
 - Consequences: Δm_{B_S} , $\sin 2\Phi_{B_S}$, wrong γ helicity
- Conclusion

Based on A. Arhrib and WSH, EPJC 27, 555 (2003)
A. Arhrib, C.K. Chua and WSH, PRD 65, 017701 (2001)
C.K. Chua, WSH and M. Nagashima, to appear

$b \rightleftharpoons s$ Phenomena: Current NP Frontier

- CP Phase $\sin 2\Phi_{B_d} \cong 0.8$ Agree with CKM Fit

but $v.$ $\begin{cases} \gamma/\phi_3 \cong 60^\circ & \text{CKM Fit} \\ \gamma/\phi_3 \sim 90^\circ & \text{B } K\pi/\pi\pi \text{ Modes} \end{cases}$

- Although $\sin 2\Phi_{B_d} \eta' K_S \cong \sin 2\Phi_{B_d}$
but $\sin 2\Phi_{B_d} \phi K_S$ **Opposite Sign**!?

- $b \rightarrow s \ell^+ \ell^- > \text{SM?}$

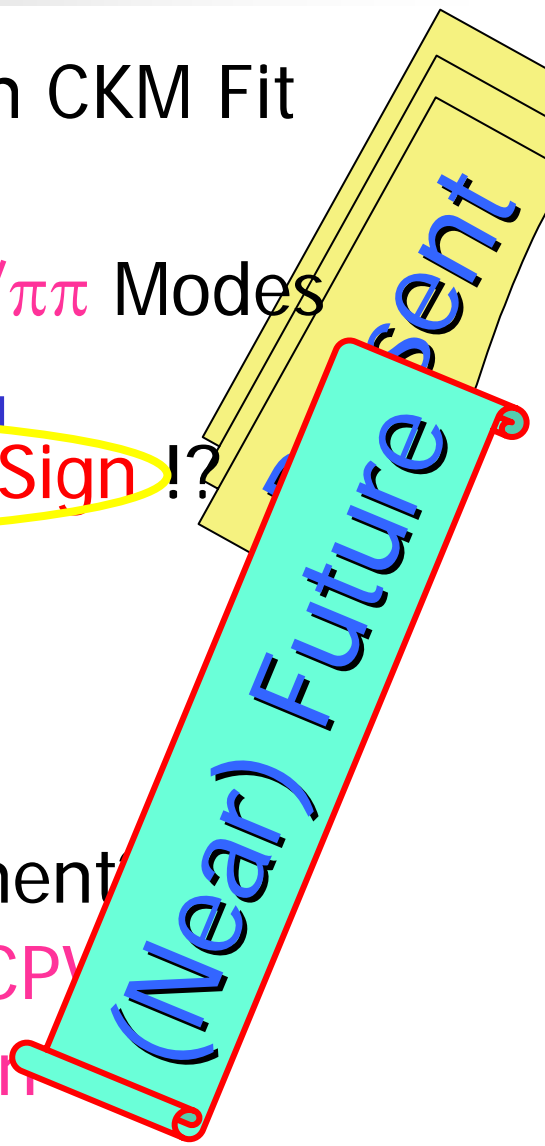
- Δm_{B_s} "Just Around the Corner"

but if $\sin 2\Phi_{B_s} \neq 0 \Rightarrow \text{NP}$

- $b \rightarrow s \gamma_L$? Wrong Helicity Component

Test via $B_s \rightarrow \phi \gamma$ Mixing-dep CPV

$\Lambda_b \rightarrow \Lambda \gamma$ Λ Polarization



Models w/ *1 Extra CP Phase* in $b \rightleftharpoons s$

- 4th Generation: In general 3 New CP Phases

Restrict to $b \rightleftharpoons s$ Only: *Just One New CP Phase*

$$\lambda_u + \lambda_c + \lambda_t + \lambda_{t'} = 0 \quad \text{Unitarity, } \lambda_i \equiv V_{is}^* V_{ib}$$

$$\approx O(\lambda^5) + \lambda^2 + \lambda_t + \lambda_{t'} \quad \Rightarrow \quad \lambda_t \approx -\lambda_c - \lambda_{t'}$$

Phase Φ_s Neglected in Literature !

$$\equiv r_s e^{i\Phi_s}$$

- Large s_R - b_R Mixing w/ SUSY: Light $\tilde{s}\tilde{b}_{1R}$ Squark

$s_R - b_R$ Mixing $\xleftrightarrow{\text{Near Maximal?}}$ $\nu_\mu - \nu_\tau$ Mixing

Flavor Symmetry
 \oplus Supersymmetry

Just One New CP Phase

in $\tilde{s}_R - \tilde{b}_R$ Mixing

➤ Flavor-driven Light s - b Flavored Squark $\tilde{s}\tilde{b}_{1R}$

4th Generation Effect

Formalism

$$\lambda_{t'} \equiv V_{t's}^* V_{t'b} \equiv r_s e^{i\Phi_s}$$

Phase Φ_s has been Neglected in Literature

$$\lambda_t \approx -\lambda_c - \lambda_{t'} \cong -\lambda^2 - \lambda_{t'} \quad \text{if } |\lambda_{t'}| \gg |\lambda_u| \approx O(\lambda^5)$$

$$\cong -0.04 \qquad \qquad \qquad \sim 0.0006$$

Take Hamiltonian for $b \quad s\gamma$

$$\mathcal{H}_{\text{eff}} = \frac{4G_F}{\sqrt{2}} \sum_{i=1}^{i=8} [\lambda_t C_i^{\text{SM}}(\mu) + \lambda_{t'} C_i^{\text{new}}(\mu)] O_i(\mu),$$

usual term

$$\lambda_t C_i^{\text{SM}} + \lambda_{t'} C_i^{\text{New}} = -\lambda_c C_i^{\text{SM}} + \lambda_{t'} (C_i^{\text{New}} - C_i^{\text{SM}})$$

Parameters

$$m_{t'}; r_s, \Phi_s$$

Satisfies GIM $\begin{cases} \lambda_{t'} \rightarrow 0 \\ m_{t'} \rightarrow m_t \end{cases}$

popular: absorb $\lambda_{t'}$ in C_i^{new} def.

Constraints: $b \quad s\gamma, \quad \Delta m_{B_S}$

Expt.

$$\left\{ \begin{array}{l} \mathcal{B}(B \rightarrow X_s \gamma) = (3.3 \pm 0.4) \times 10^{-4} \\ \Delta m_{B_S} > 14.9 \text{ ps}^{-1} \end{array} \right.$$

Choose param.

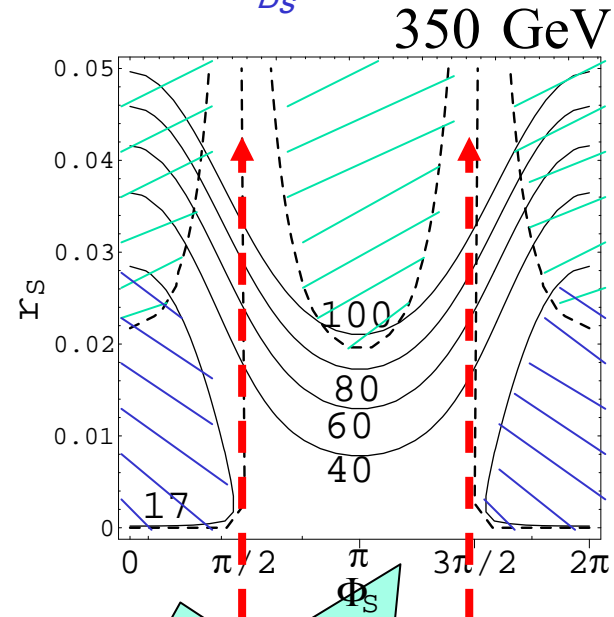
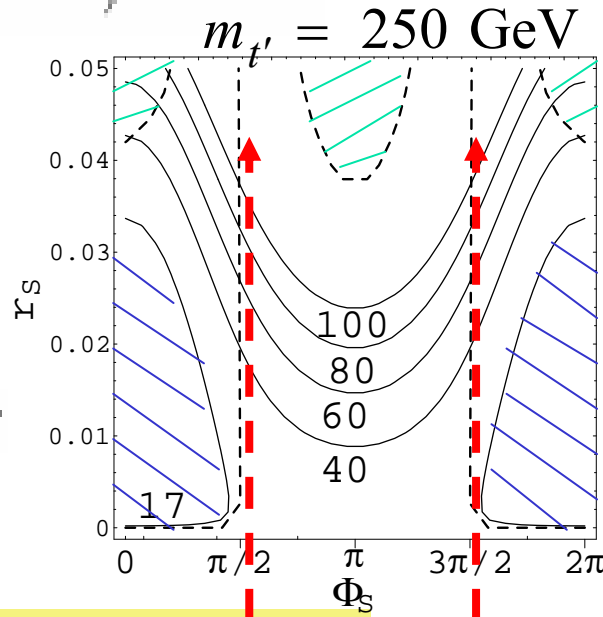
\Rightarrow SM Values

$$\mathcal{B}(B \rightarrow X_s \gamma) \simeq 3.26 \times 10^{-4}$$

$$\Delta m_{B_S}^{\text{SM}} \simeq 17.0 \text{ ps}^{-1}$$

dashed: $b \quad s\gamma,$

solid: Δm_{B_S}



- $b \quad s\gamma, \quad \Delta m_{B_S}$ Complementary
 - $b \quad s\gamma$ Weak $m_{t'}$ -dep
 - \Rightarrow Not So Constraining
 - Δm_{B_S} Strong $m_{t'}$ -dep
 - \Rightarrow Not Constructive Interf.

$\Phi_S = \pi/2, 3\pi/2$
Most Forgiving
 \therefore Add in quadrature

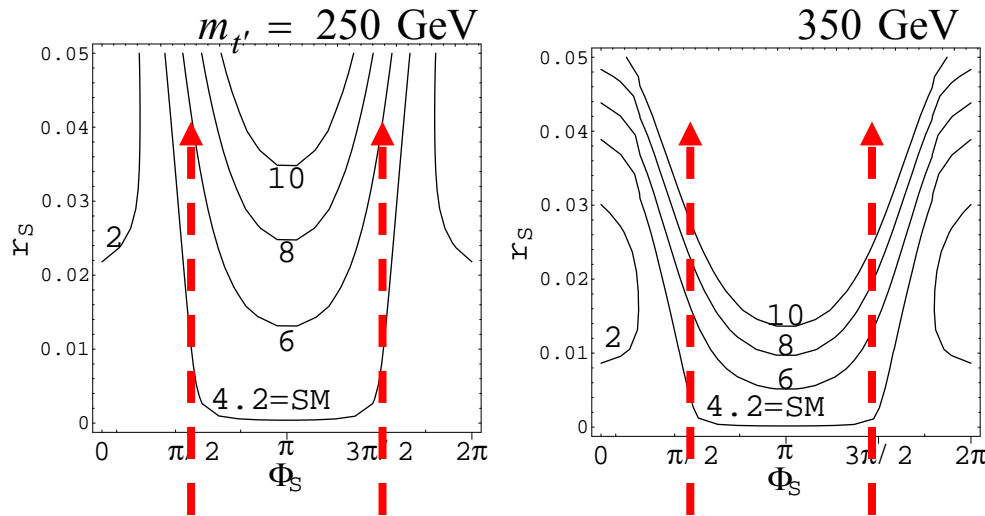
Constraint/Test: $b \rightarrow s \ell^+ \ell^-$ Rate

EW Penguin has Strong $m_{t'}$ -dep., Just Like B_s Mixing

$$\left. \begin{aligned} \mathcal{B}(B \rightarrow K \ell^+ \ell^-) &\sim (0.5-0.8) \times 10^{-6} \\ \mathcal{B}(B \rightarrow K^* \ell^+ \ell^-) &\sim (0.8-1.7) \times 10^{-6} \end{aligned} \right\} \text{Yet to Settle (a bit high)}$$

$$\mathcal{B}(B \rightarrow X_s \ell^+ \ell^-) = (6.1 \pm 1.4^{+1.3}_{-1.1}) \times 10^{-6}$$

Less hadronic uncertainty
Maybe high vs. NNLO



SM ~ NNLO SM3

- More Stringent than Δm_{B_s}
 \Rightarrow *Improve Soon!*
- $\Phi_s \sim \pi/2, 3\pi/2$
 Allow Larger r_s
- May be Needed If
 NNLO Remains Low
- No Further Info in $m^2(\ell\ell)$
- Constrained by BR
 $\Rightarrow A_{FB} \sim \text{SM}$

The Highlight: Large $\sin 2\Phi_{B_S} \neq 0$

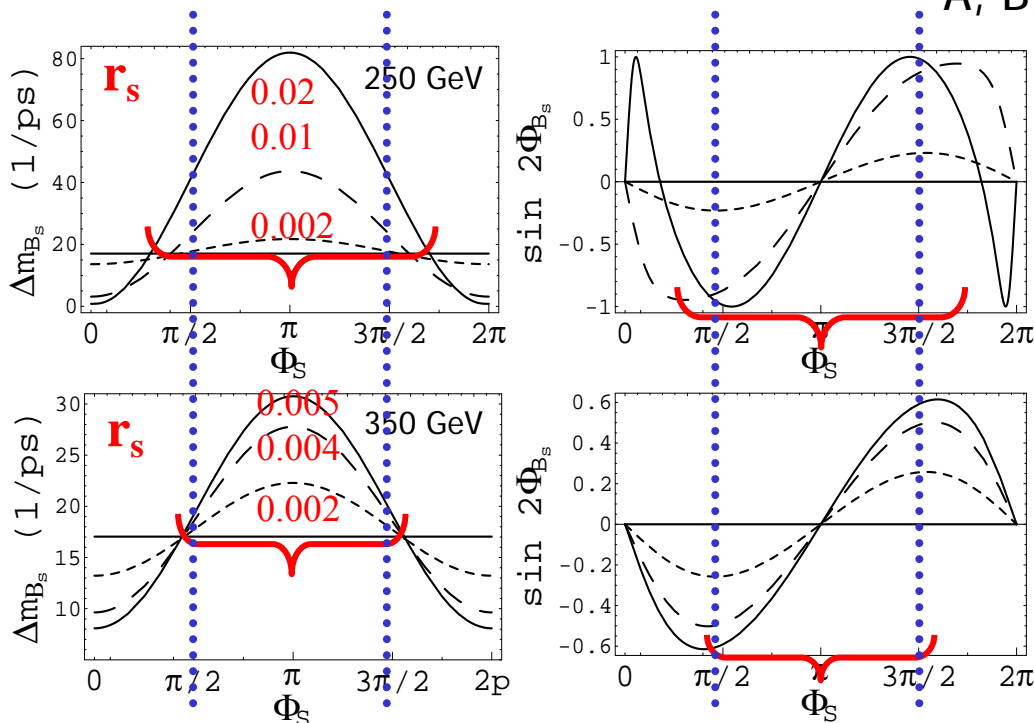
- $\sin 2\Phi_{B_d}$ Agree with SM3, but May be Affected by $V_{t'd}^* V_{t'b}$
- $\sin 2\Phi_{B_S} = 0$ Strictly in SM3 \Rightarrow Any $\sin 2\Phi_{B_S} \neq 0$ Implies NP !!

Def.

$$M_{12} = |M_{12}| e^{2i\Phi_{B_S}} \approx r_s^2 e^{2i\Phi_S} A + r_s e^{i\Phi_S} B + C$$

$$\Delta m_{B_S} = 2|M_{12}|$$

A, B, C are quadratic fns of m_t, m_t



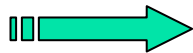
- Even for Small r_s
 $\sin 2\Phi_{B_S}$ Can be Large
- Both Signs Possible
- Largest If Δm_{B_S}
"Just Around Corner"
i.e. when $\Phi_S \sim \pi/2, 3\pi/2$
 $\left\{ \begin{array}{l} \Delta m_{B_S} \text{ Quite Measurable} \\ \sin 2\Phi_{B_S} \sim \pm 1 \end{array} \right.$

Good at Tevatron !

But, ...

4th Generation Not Effective on EM & Strong Penguins

- $b \rightarrow s\gamma$ affected only when r_s (and $m_{t'}$) gets large
No Mixing CPV because (L) Chirality Same as SM
- $b \rightarrow sg^*$ Even Less Affected
 \Rightarrow Cannot Touch $S_{\phi K_S}, S_{\eta' K_S}$ etc.



Turn to one that Can

Light

$$\tilde{s}b_{1R}$$

& $S_{\phi K_S}, S_{\eta' K_S}, \Delta m_{B_S}, \sin 2\Phi_{B_S}$

Mass/Mixing Hierarchy

$$\hat{M}_u = \frac{M_u}{m_t} \sim \begin{bmatrix} \lambda^7 & \lambda^5 & \lambda^3 \\ ? & \lambda^4 & \lambda^2 \\ ? & ? & 1 \end{bmatrix}, \quad \hat{M}_d = \frac{M_d}{m_b} \sim \begin{bmatrix} \lambda^4 & \lambda^3 & \lambda^3 \\ ? & \lambda^2 & \lambda^2 \\ ? & ? & 1 \end{bmatrix}$$

Abelian Flavor Symmetry

$$M_{ij}M_{ji} \approx M_{ii}M_{jj} \quad \text{Commuting Charges}$$

$$\hat{M}_u \sim \begin{bmatrix} \lambda^7 & \lambda^5 & \lambda^3 \\ \lambda^6 & \lambda^4 & \lambda^2 \\ \lambda^4 & \lambda^2 & 1 \end{bmatrix}, \quad \hat{M}_d \sim \begin{bmatrix} \lambda^4 & [\lambda^3] & [\lambda^3] \\ [\lambda^3] & \lambda^2 & \lambda^2 \\ [\lambda] & 1 & 1 \end{bmatrix}$$

Prominent

Nir-Seiberg, PLB'93; Leurer-Nir-Seiberg, NPB'94

but

N.B. Need 4 Texture Zeros (Chua-WSH, PRL'01) : Set [..] = 0 right-handed

Arhrib-Chua-WSH, PRD'01 *Decouple d flavor*

Alternative Picture:

$$\nu_\mu - \nu_\tau \text{ Mixing} \xrightarrow{\text{Near Maximal?}} s_R - b_R \text{ Mixing}$$

Chang,
Masiero,
Murayama

Right-handed Quarks **Inert in SM**

→ Right-handed **S**quarks

$\tilde{d}_{iR} - d_{jR} - \tilde{g}$ Dynamics



SUSY

Assume

Squarks $\approx \tilde{m}$, almost degenerate
AFS not far above SUSY Scale

$$\left(\tilde{M}_q^2\right)_{RL} = \left(\tilde{M}_q^2\right)_{LR}^{*T} \approx M_q \tilde{m},$$

also important
 $\frac{m_q}{\tilde{m}}$ suppressed

$$\left(\tilde{M}_Q^2\right)_{LL} \approx \tilde{m}^2 V_{\text{CKM}},$$

CKM suppressed

but $\left(\tilde{M}_d^2\right)_{RR} \approx \tilde{m}^2$

$$\begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 1 \\ 0 & 1 & 1 \end{bmatrix}$$

RR Sector
Impact $b \leftrightarrow s$
thru **SUSY**

Level Splitting by Large Mixing

→ Drive One State Light

$$(\tilde{M}_q^2)_{RR} \approx \tilde{m}^2 \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 1 \\ 0 & 1 & 1 \end{bmatrix} \Rightarrow \tilde{M}_{RR}^{2(sb)} = \begin{bmatrix} \tilde{m}_{22}^2 & \tilde{m}_{23}^2 e^{-i\sigma} \\ \tilde{m}_{23}^2 e^{i\sigma} & \tilde{m}_{33}^2 \end{bmatrix} = R \begin{bmatrix} \tilde{m}_1^2 & 0 \\ 0 & \tilde{m}_2^2 \end{bmatrix} R^\dagger$$

$$R = \begin{bmatrix} c_\theta & s_\theta \\ -s_\theta e^{i\sigma} & c_\theta e^{i\sigma} \end{bmatrix}$$

- Near Maximal Mixing: $s_\theta \sim 1/\sqrt{2}$
- $\sigma = \arg(\tilde{M}_d^2)_{RR}^{32}$ similar footing as $\phi_3 \equiv \arg V_{ub}^*$

1 CP Phase

Light $\bar{s}b_1$ w/ Some Tuning:

set $\tilde{m}_{22}^2 = \tilde{m}_{33}^2 = \tilde{m}^2$,

With $\tilde{m}_{23}^2/\tilde{m}^2 \equiv 1 - \delta \simeq 1$,

$$\tilde{m}_1^2 \cong \delta \tilde{m}^2, \quad \tilde{m}_2^2 \cong (2 - \delta) \tilde{m}^2$$

Tunings:

take $\tilde{m} = 2$ (1) TeV

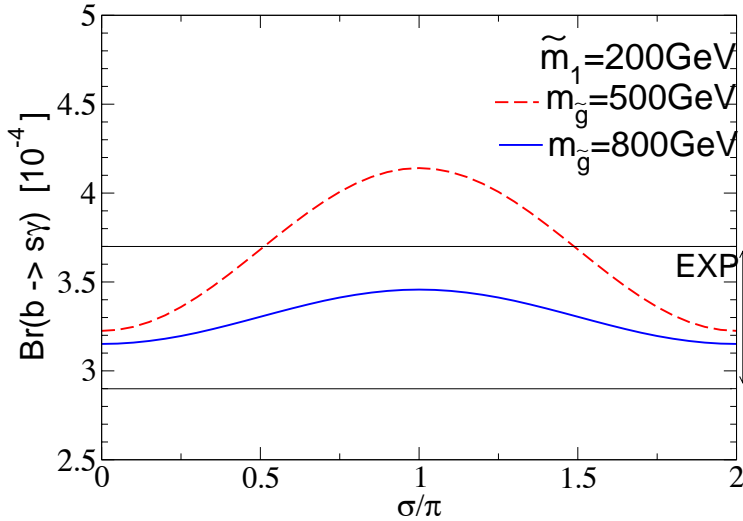
$$\delta = \lambda^2, \quad \lambda^3, \quad \lambda^4 \quad (\lambda, \quad \lambda^2, \quad \lambda^3)$$

$$\Rightarrow \tilde{m}_1 = 440, \quad 206, \quad 97 \text{ GeV} \quad (470, \quad 220, \quad 103 \text{ GeV})$$

Focus on
200s GeV

Rates

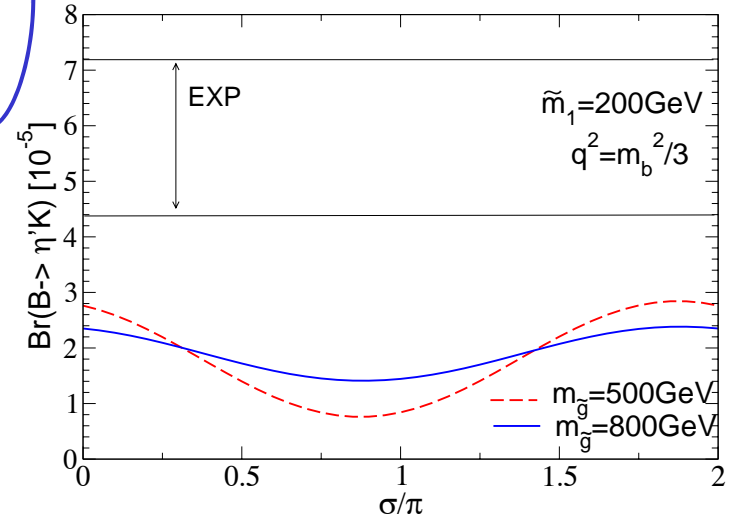
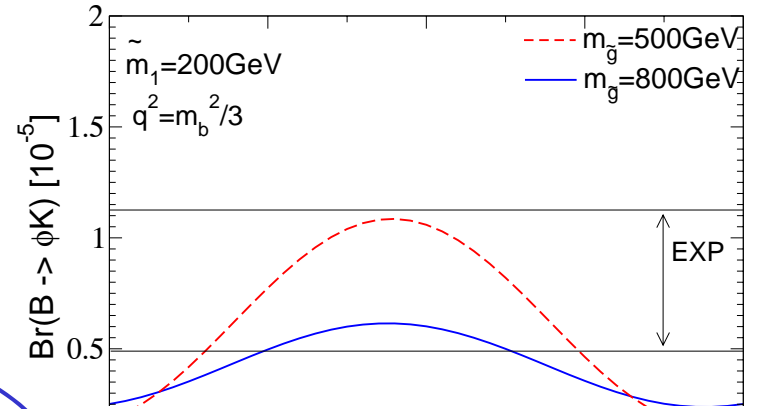
- Survive $b \rightarrow s \gamma$ Constraint



$\tilde{m} = 2 \text{ TeV} \Rightarrow m_{\tilde{g}} = 800 \text{ GeV}$
 More Conservative

Cannot Account of $\eta' K_s$ Rate
 ϕK_s Rate also a bit Low
 But, Not New Problems

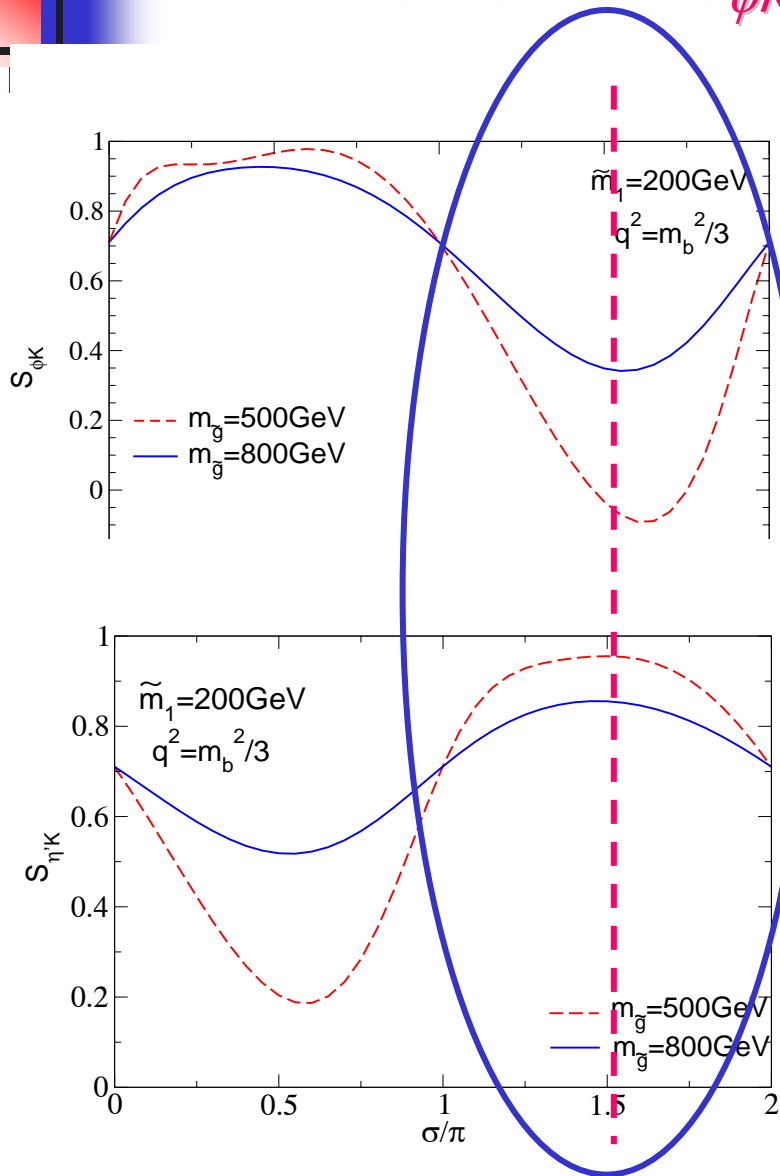
- $B \rightarrow \phi K_s, \eta' K_s$ Rates



A Little Note on Formalism ...

- Besides $O_{1,2}$ Tree
 - O_{3-6} Strong Penguin
 - O_{7-10} EM/EW Penguin
 - also* $O_{11,12}$ γ/g Dipole
- Right-handed Coupling $\rightarrow O'_i$

Tune for $S_{\phi K_S}$, $S_{\eta' K_S}$ Possible



- $S_{\phi K_S} \lesssim 0$ prefers lower σ -plane

$\Rightarrow S_{\eta' K_S} \cong \sin 2\Phi_{B_d}$ as well !

- Lower gluino mass lowers $S_{\phi K_S}$
- Hadronic Parameters

CP phase enters $B \rightarrow \phi K_S$ via

$$\frac{\alpha_s}{4\pi} \frac{m_b^2}{q^2} (C_{12} + C'_{12}) \bar{S}_{\phi K_S}$$

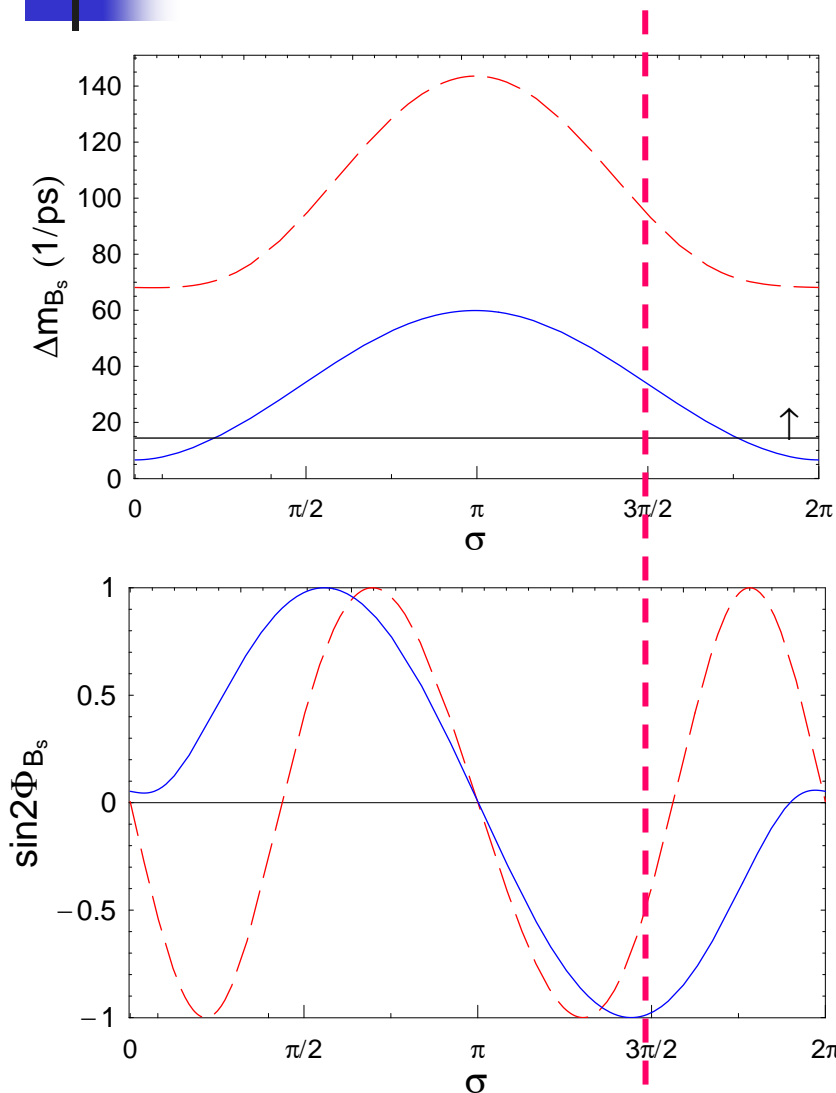
Lower q^2 , larger $\bar{S}_{\phi K_S}$

\Rightarrow Larger $|S_{\phi K_S}|$

- If $B \rightarrow \eta' K_S$ rate due to (SM) QCD Effect $\Rightarrow S_{\eta' K_S}$ drops a bit
- Could $\sigma \approx 3\pi/2$ be it?

noticed also by Khalil & Kou
Murayama et al.

Consequences: Δm_{B_S} , $\sin 2\Phi_{B_S}$



Take $\sigma \approx 3\pi/2$ as example

- 800 GeV gluino “preferred”
 - $\Delta m_{B_S} \sim 20\text{-}40 \text{ ps}^{-1}$ “Easy”
 - $\sin 2\Phi_{B_S} \sim -1$ Super
- Lighter gluino at 500 GeV
 - $\Delta m_{B_S} > 60 \text{ ps}^{-1}$ Tougher
 - $\sin 2\Phi_{B_S} \sim -1 - 0$
- For lighter gluino, periodicity change

$$\therefore M_{12} \equiv |M_{12}| e^{2i\Phi_{B_S}} \equiv a e^{-2i\sigma} + b e^{-i\sigma} + c$$

a -term dominant (two $\widetilde{s}b_{1R}$ exch.)

- 100 GeV $\widetilde{s}b_{1R} \sim 200$ GeV case
Except, **Direct Detection !**

More: Wrong Helicity Photon in $b \rightarrow s \gamma_L$

- Strength of Mixing-dep. CPV

$$\sin 2\theta = \frac{2 |C_7 C'_7|}{(|C_7|^2 + |C'_7|^2)}$$

in, e.g. $B_s \rightarrow \phi \gamma$

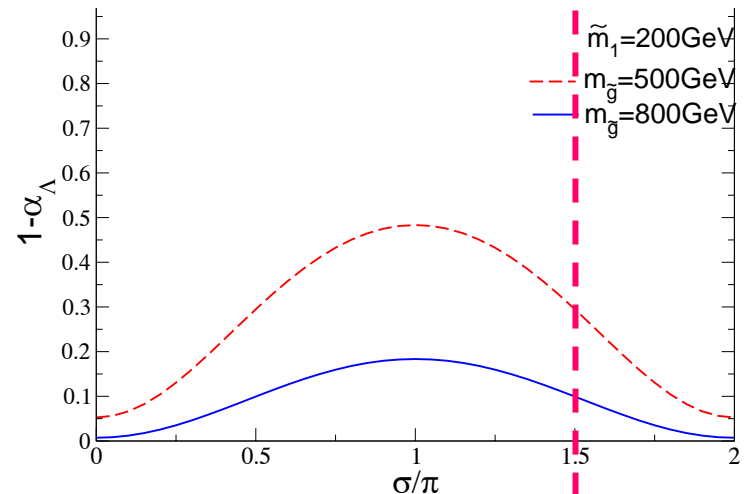
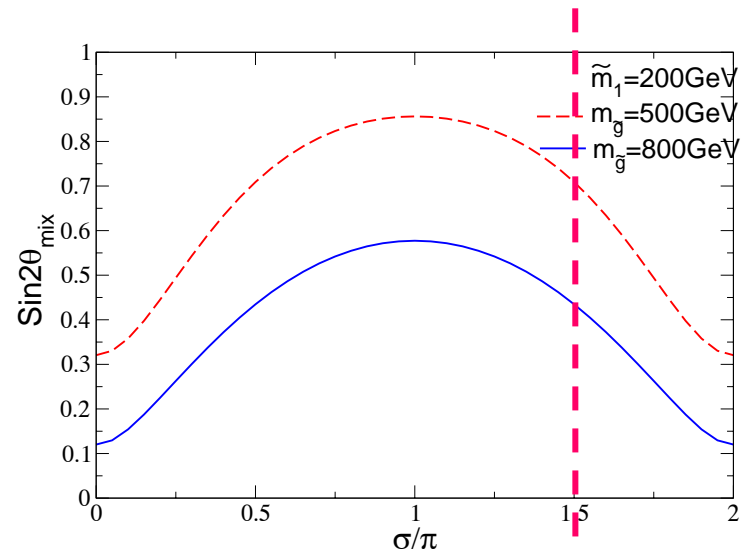
Test γ_L component

(Atwood, Gronau, Soni, 1997)

- Can also test γ_L comp. via Δ Polarization in $\Lambda_b \rightarrow \Lambda \gamma$

(Mannel, Recksiegel, 1997)

All Testable at Tevatron and Beyond



Conclusion

- Impact of 4th Generation (Emphasized **New CP Phase**)
Prominent in - **B_s Mixing & CPV**
- **EW $b \rightarrow s$ Penguin**

Mild in EM/Strong Penguin

- A Light Flavor-mixed $\tilde{s}b_{1R}$ Squark?
 - Survive $b \rightarrow s\gamma$
 - Can Account for $S_{\phi K_S}, S_{\eta' K_S}$
 - Strong Impact on: $\Delta m_{B_S}, \sin 2\Phi_{B_S}, b \rightarrow s\gamma_L$

⇒ **Study at B Factories, Tevatron, and ...**